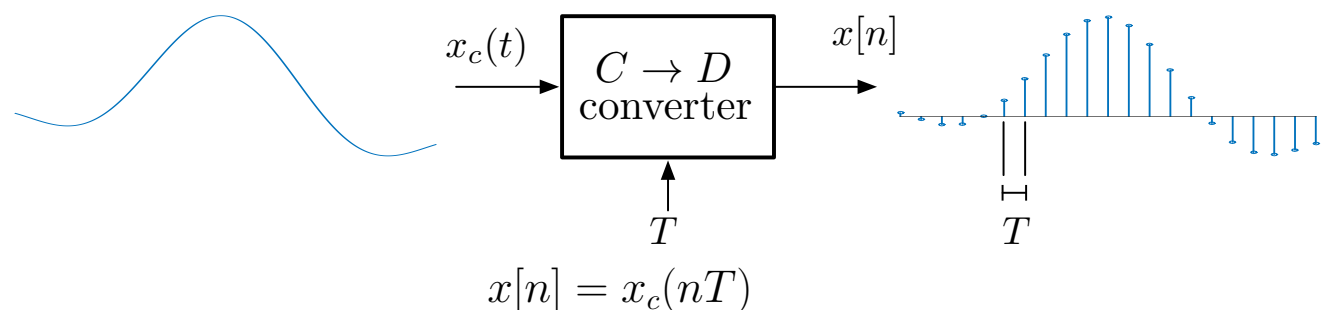


# I. Signal Discretization using Basis Decompositions

We will start by reviewing one of the foundational results of digital signal processing: the Shannon-Nyquist sampling theorem. We will use this result as a first example of how continuous-time signals can be systematically discretized (translated into a discrete list of numbers).

## The Shannon-Nyquist sampling theorem

*Sampling* turns a continuous-time signal  $x_c(t)$  into a discrete list of numbers simply by evaluating it at equally spaced points:



( $C \rightarrow D$ =continuous-to-discrete.)

The constant  $T$  is the **sampling interval** (the amount of time that passes between each sample).

This is a very common practice, and there exists very sophisticated hardware that implements it. Examples:

- Texas Instruments makes an ADC, the 12J4000, that takes 4 *billion* samples per second ( $T = 0.25$  nanoseconds) at a (reported) resolution of 12 bits. Cost:  $\approx$ \$2000.
- Another ADC from TI, the TIADS1278, takes 144,000 samples per second ( $T = 6.94$  microseconds) at a (reported) resolution of 24 bits. Cost:  $\approx$ \$25

## Questions:

1. When can you reconstruct  $x_c(t)$  perfectly from its samples?
2. How do you do it?

## Answers:

1. When  $x_c(t)$  is **bandlimited**, i. e. when

$$X_c(j\Omega) = 0 \quad \text{for all} \quad |\Omega| > \pi/T$$

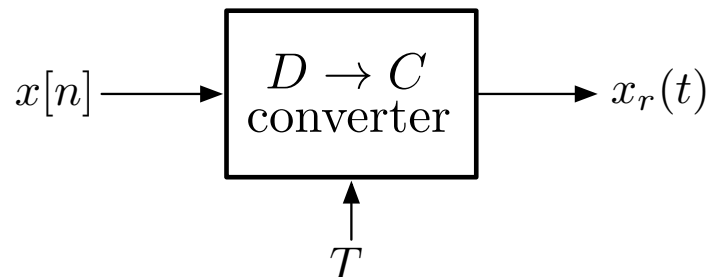
where  $X_c(j\Omega)$  is the continuous time Fourier transform (CTFT) of  $x_c(t)$ :

$$X_c(j\Omega) = \int_{-\infty}^{\infty} x_c(t)e^{-j\Omega t} dt, \quad \text{or} \quad \mathbf{X}_c = \mathcal{F}\{\mathbf{x}_c\}$$

In other words, the sampling rate ( $= 1/T$  in Hz, or  $2\pi/T$  in rad/sec) must be at least **twice** the maximum frequency present in the signal.

This is known as the **Nyquist criterion**.

2. We reconstruct the continuous time signal from the discrete sample sequence using **sinc interpolation**:



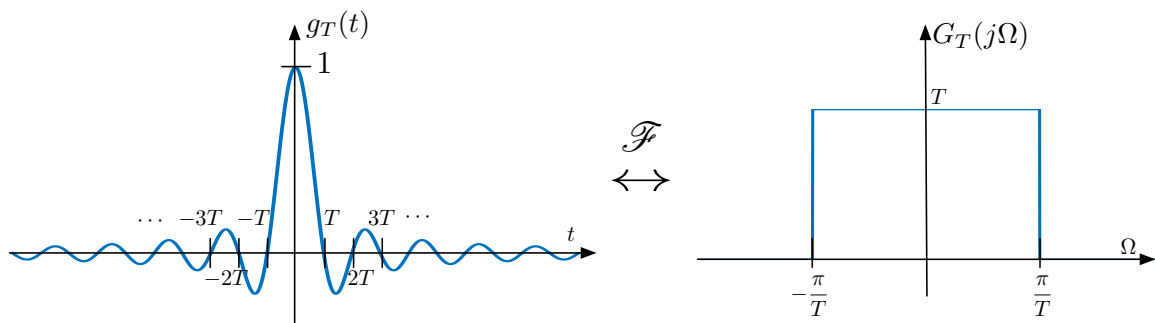
( $D \rightarrow C$  = discrete-to-continuous.)

Mathematically, we can write the output as:

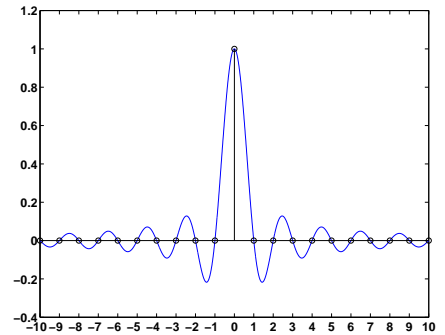
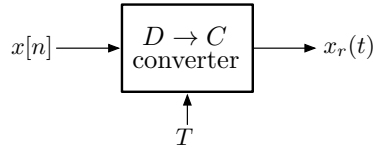
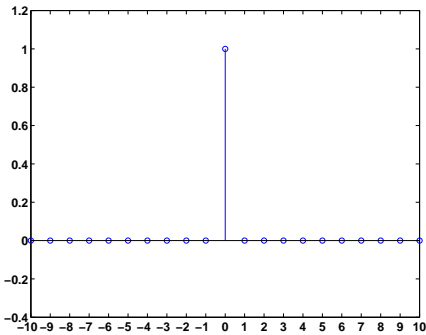
$$\begin{aligned}
 x_r(t) &= \sum_{n=-\infty}^{\infty} x[n] \frac{\sin(\pi(t - nT)/T)}{\pi(t - nT)/T} \\
 &= \sum_{n=-\infty}^{\infty} x[n] \underbrace{g_T(t - nT)}_{\text{shifts of the sinc}}
 \end{aligned}$$

Recall that:

$$g_T(t) = \frac{\sin(\pi t/T)}{\pi t/T} \quad \xleftrightarrow{\mathcal{F}} \quad G_T(j\Omega) = \begin{cases} T, & |\Omega| \leq \frac{\pi}{T}, \\ 0, & |\Omega| > \frac{\pi}{T} \end{cases}$$

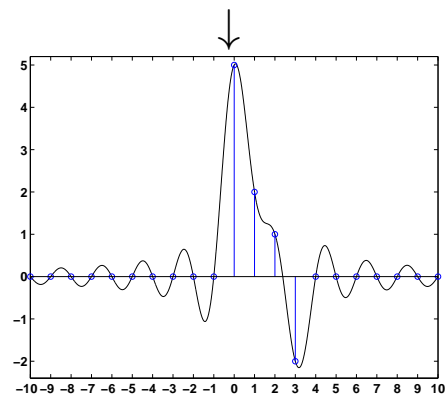
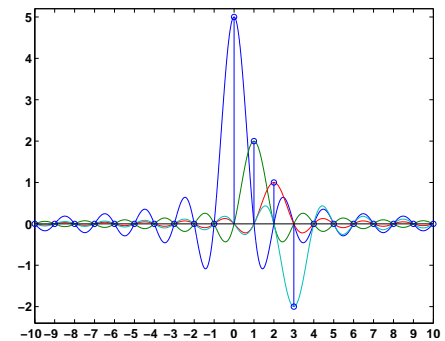
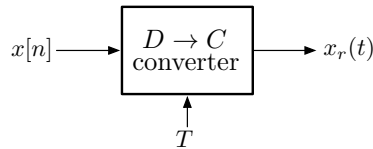
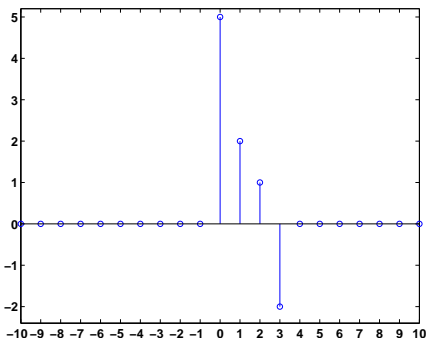


Single sample:



Notice that the sinc function is exactly zero at the other sample locations.

Multiple samples:



In between samples, multiple sincs combine to yield a smooth signal.

## The Fundamental Theorem of DSP

If  $x_c(t)$  is bandlimited to  $B$  ( $X_c(j\Omega) = 0$  for  $|\Omega| > B$ ), then it can be perfectly reconstructed from samples spaced  $T \leq \pi/B$  apart:

$$x_c(t) = \sum_{n=-\infty}^{\infty} x[n] g_T(t - nT),$$

where

$$x[n] = x_c(nT), \quad g_T(t) = \frac{\sin(\pi t/T)}{\pi t/T}.$$

1. This is known as the **Shannon-Nyquist sampling theorem**
2. It is the backbone of DSP — it essentially says that we can process  $x_c(t)$  by processing its samples
3. The samples are a discrete list of numbers, and hence can be processed **digitally** on a computer, giving us tremendous flexibility.
4. The two equations above are our first example of a **reproducing formula**, which shows how a signal can be written as a discrete combination of linear functionals of that signal (samples, in this case) weighted against a set of fixed “basis” signals. This is a central theme in this first section of the course.

## Frequency domain interpretation

Like many things, it is illuminating to look at sampling and reconstruction in the frequency domain.

First, we will relate the **discrete time Fourier transform** (DTFT) of  $x[n]$  to the CTFT of  $x_c(t)$ :

$$\begin{aligned} X(e^{j\omega}) &= \sum_{n=-\infty}^{\infty} x[n] e^{-j\omega n} = \sum_{n=-\infty}^{\infty} x_c(nT) e^{-j\omega n} \\ &= \sum_{n=-\infty}^{\infty} \left( \frac{1}{2\pi} \int_{-\infty}^{\infty} X_c(j\Omega) e^{j\Omega nT} d\Omega \right) e^{-j\omega n} \\ &= \frac{1}{2\pi} \int_{-\infty}^{\infty} X_c(j\Omega) \cdot \left( \sum_{n=-\infty}^{\infty} e^{jn(\Omega T - \omega)} \right) d\Omega. \end{aligned}$$

Recall the **Poisson Summation Formula**:

$$\sum_{n=-\infty}^{\infty} e^{jn\omega} = 2\pi \sum_{k=-\infty}^{\infty} \delta(\omega - 2\pi k)$$

where

$$\delta(\omega) = \text{“Dirac delta function”}.$$

Plugging this in, we have

$$\begin{aligned} X(e^{j\omega}) &= \int_{-\infty}^{\infty} X_c(j\Omega) \cdot \sum_{k=-\infty}^{\infty} \delta(\Omega T - \omega - 2\pi k) d\Omega \\ &= \sum_{k=-\infty}^{\infty} \int_{-\infty}^{\infty} X_c(j\Omega) \delta(\Omega T - \omega - 2\pi k) d\Omega \\ &= \frac{1}{T} \sum_{k=-\infty}^{\infty} X_c \left( j \left( \frac{\omega + 2\pi k}{T} \right) \right) \end{aligned}$$

There are essentially two things going on here:

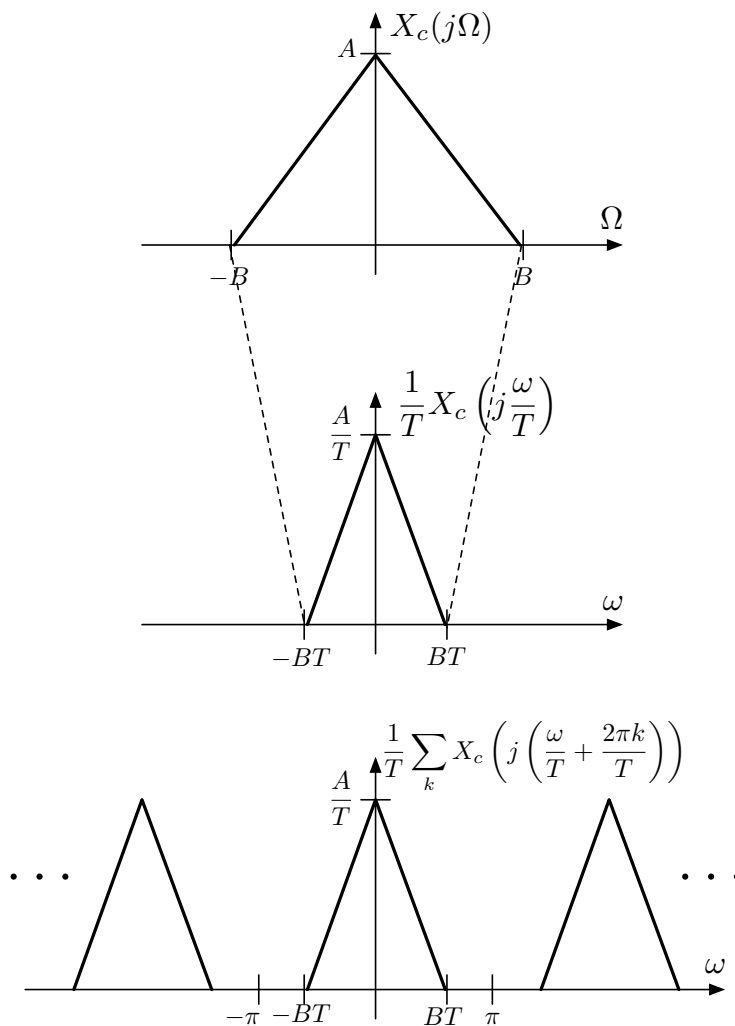
1.  $X_c(j\Omega) \longrightarrow \frac{1}{T}X_c\left(j\frac{\omega}{T}\right)$

**dilates** the spectrum

2.  $\frac{1}{T}X_c\left(j\frac{\omega}{T}\right) \longrightarrow \frac{1}{T}\sum_k X_c\left(j\left(\frac{\omega}{T} + \frac{2\pi k}{T}\right)\right)$

makes this dilation **periodic** (w/ period  $2\pi$ ).

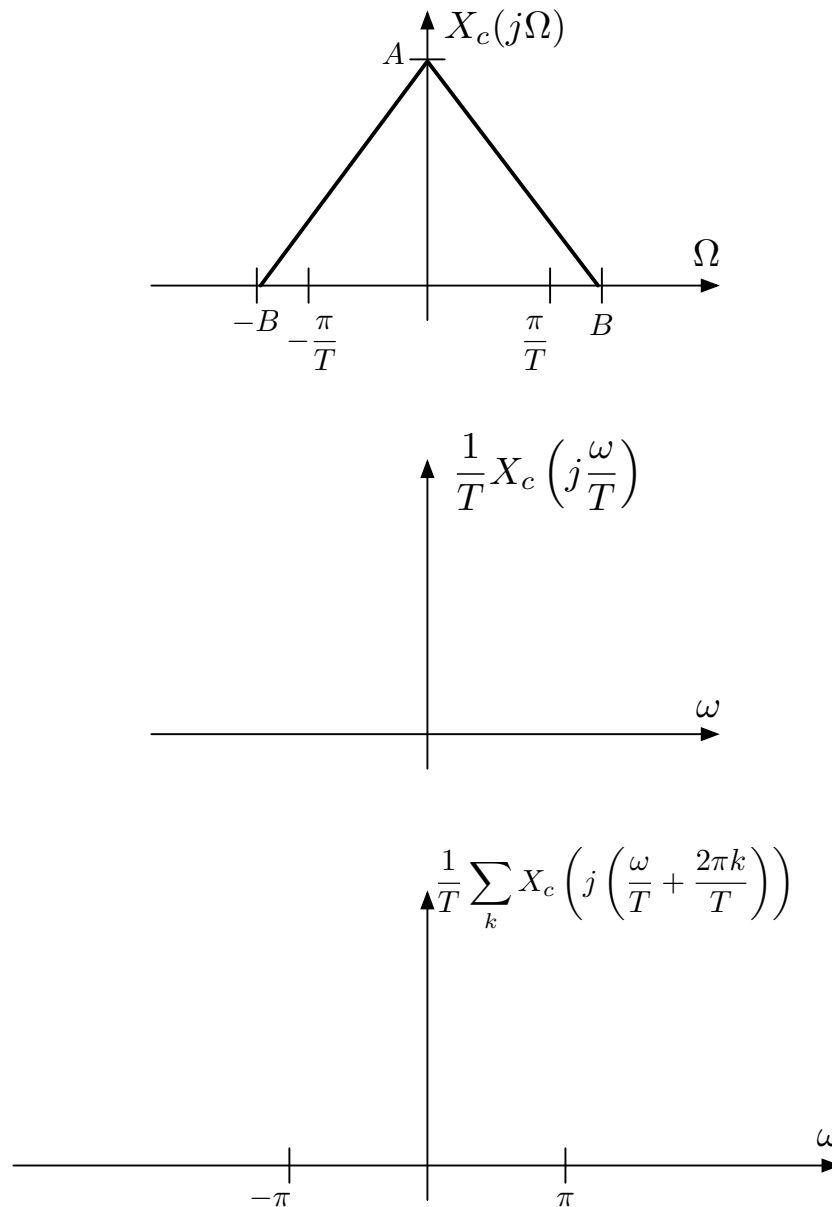
Graphically, this is what happens for  $B < \pi/T$ :





## Aliasing

If  $T > \pi/B$ , there is trouble:



What is another signal with the same samples as  $x_c(t)$ ?

## Reconstruction

The reconstructed signal is

$$x_r(t) = \sum_{n=-\infty}^{\infty} x[n] g_T(t - nT),$$

and so

$$\begin{aligned} X_r(j\Omega) &= \sum_{n=-\infty}^{\infty} x[n] G_T(j\Omega) e^{-j\Omega nT} \\ &= G_T(j\Omega) \cdot \sum_{n=-\infty}^{\infty} x[n] e^{-j\Omega nT} \\ &= G_T(j\Omega) X(e^{j\Omega T}) \end{aligned}$$

Again, there are two steps:

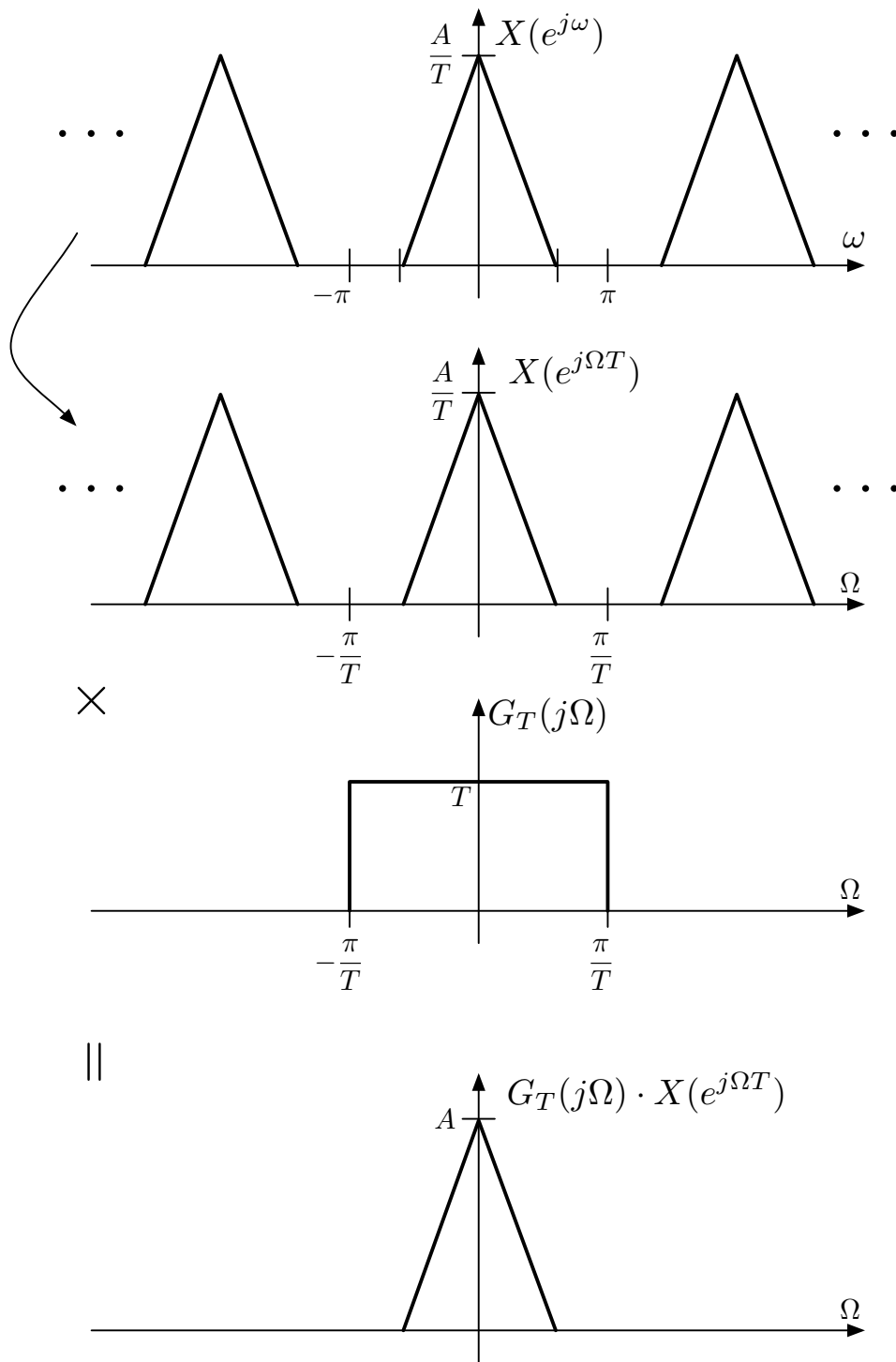
1.  $X(e^{j\omega}) \longrightarrow X(e^{j\Omega T})$

**dilates** the (periodic) spectrum

2.  $X(e^{j\Omega T}) \longrightarrow G_T(j\Omega) \cdot X(e^{j\Omega T})$

restricts the spectrum to its fundamental period

Graphically, this is what happens:

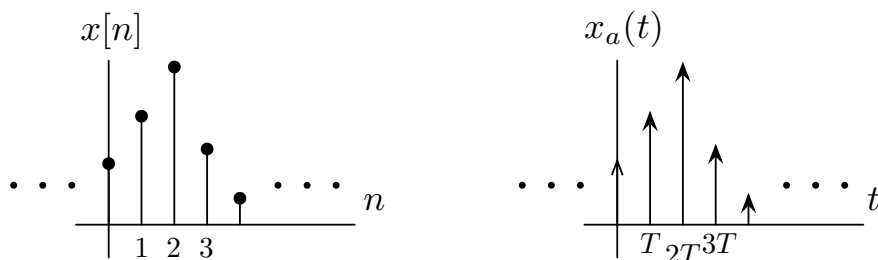


A little more on the  $X(e^{j\omega}) \longrightarrow X(e^{j\Omega T})$  step ...

What we are doing is taking a discrete sequence  $x[n]$  (with DTFT  $X(e^{j\omega})$ ) and turning it into a function  $x_a(t)$  (with CTFT  $X_a(j\Omega) = X(e^{j\Omega T})$ ) of a continuous time variable.

Set

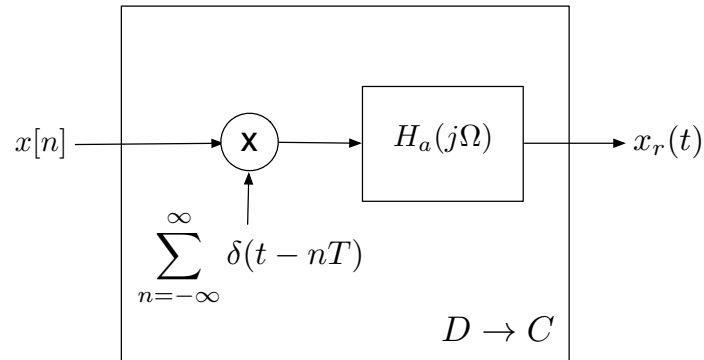
$$x_a(t) = \sum_{n=-\infty}^{\infty} x[n] \delta(t - nT)$$



Then

$$\begin{aligned} X_a(j\Omega) &= \int_{-\infty}^{\infty} \sum_n x[n] \delta(t - nT) e^{-j\Omega t} dt \\ &= \sum_n x[n] \int_{-\infty}^{\infty} \delta(t - nT) e^{-j\Omega t} dt \\ &= \sum_n x[n] e^{-j\Omega T n} = X(e^{j\Omega T}) \end{aligned}$$

So the  $D \rightarrow C$  converter converts the sample sequence into a **spike train** and then low pass filters it. We can interpret what is inside this block as:

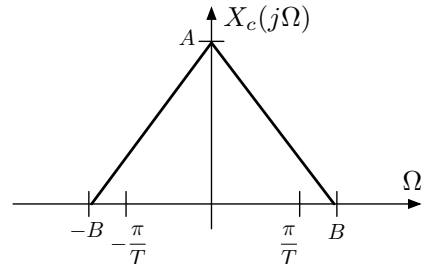


where

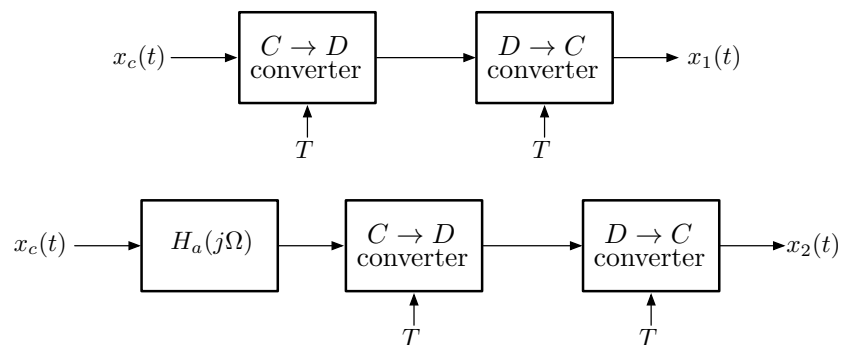
$$H_a(j\Omega) = \begin{cases} T, & |\Omega| \leq \frac{\pi}{T}, \\ 0, & \text{otherwise.} \end{cases}$$

## Anti-aliasing filters

Suppose the spectrum of  $x_c(t)$  looks like



Compare the outputs of these two systems:



where

$$H_a(j\Omega) = \begin{cases} 1 & |\Omega| \leq \pi/T \\ 0 & |\Omega| > \pi/T \end{cases}.$$

Which is closer to  $x_c(t)$ ?

That is, which is smaller:

$$\int |x_c(t) - x_1(t)|^2 dt \quad \text{or} \quad \int |x_c(t) - x_2(t)|^2 dt \quad ?$$

# Appendix: Technical Review

## The Dirac delta function

The Dirac delta is a *generalized function*, defined through the relation

$$\int_{-L}^L x(t)\delta(t) dt = x(0), \quad \text{for any } L > 0.$$

More generally,

$$\int_{t \in \mathcal{T}} x(t)\delta(t - t_0) dt = \begin{cases} x(t_0), & \text{if } t_0 \in \mathcal{T}, \\ 0, & \text{otherwise.} \end{cases}$$

$\delta(t)$  is not a function in the usual sense, but we can manipulate algebraically in much the same way we manipulate standard functions.

The delta function is the “derivative” of the Heaviside step function

$$\mu(t) = \begin{cases} 1, & t \geq 0, \\ 0, & t < 0. \end{cases},$$

in that they obey a relation of the same form as the Fundamental Theorem of Calculus:

$$\mu(t) = \int_{-\infty}^t \delta(\tau) d\tau.$$

The formalism for  $\delta(t)$  and other generalized functions is found in the mathematical theory of distributions. A nice overview of this theory can be found in the classic text *Distributions, Complex Variables, and Fourier Transforms*, by Hans Bremermann (1965).

## The continuous-time Fourier transform (CTFT)

The CTFT of a signal  $x(t)$  is

$$X(j\Omega) = \int_{-\infty}^{\infty} x(t) e^{-j\Omega t} dt.$$

The convention of using  $j\Omega$  as the argument (instead of just  $\Omega$ ) is historical, and is common in the signal processing literature.

Anytime you see an integral expression like the one above, it is fair to ask whether or not it converges. If  $x(t)$  is *absolutely integrable*, in that

$$\int_{-\infty}^{\infty} |x(t)| dt < \infty,$$

then  $X(j\Omega)$  is well-defined for all  $\Omega \in \mathbb{R}$ . It is also bounded, as in this case

$$|X(j\Omega)| = \left| \int_{-\infty}^{\infty} x(t) e^{-j\Omega t} dt \right| \leq \int_{-\infty}^{\infty} |x(t)| |e^{-j\Omega t}| dt = \int_{-\infty}^{\infty} |x(t)| dt.$$

If  $x(t)$  has *finite energy*, in that

$$\int_{-\infty}^{\infty} |x(t)|^2 dt \leq \infty,$$

then the Fourier transform is also well-defined, but you have to be a little more careful about what it means for two functions to be equal to one another. We will talk a little more about this later, but it is really just a mathematical detail which ends up not affecting our outlook on this topic at all.

The inverse CTFT is

$$x(t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} X(j\Omega) e^{j\Omega t} d\Omega.$$



The Parseval theorem states that the energy in the time- and frequency-domains are equal to one another (to within a constant of  $1/2\pi$ ):

$$\int_{-\infty}^{\infty} |x(t)|^2 dt = \frac{1}{2\pi} \int_{-\infty}^{\infty} |X(j\Omega)|^2 d\Omega.$$

## The discrete-time Fourier transform (DTFT)

The DTFT of the sequence of numbers  $\{x[n], n \in \mathbb{Z}\}$  is defined as

$$X(e^{j\omega}) = \sum_{n=-\infty}^{\infty} x[n]e^{-j\omega n}.$$

Again, this sum is clearly well-defined (and bounded) when  $x[n]$  is absolutely summable,

$$\sum_{n=-\infty}^{\infty} |x[n]| < \infty,$$

and we can make sense of it when  $x[n]$  has finite energy,

$$\sum_{n=-\infty}^{\infty} |x[n]|^2 < \infty.$$

Notice that  $X(e^{j\omega})$  is  $2\pi$ -periodic, as

$$e^{-j\omega n} = e^{-j(\omega+2\pi\ell)n} \quad \text{for all } \ell \in \mathbb{Z}.$$

The inverse DTFT is

$$x[n] = \frac{1}{2\pi} \int_{-\pi}^{\pi} X(e^{j\omega})e^{j\omega n} d\omega.$$

The DTFT also preserves energy up to a constant, as

$$\sum_{n=-\infty}^{\infty} |x[n]|^2 = \frac{1}{2\pi} \int_{-\pi}^{\pi} |X(e^{j\omega})|^2 d\omega.$$